

Appendices for:
Discontinuous mechanical behaviors of existing shield tunnel with stiffness
reduction at longitudinal joints

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Appendix A. List of symbols

A : cross-sectional area
 D : diameter of existing shield tunnel/pipeline
 EI : bending stiffness of beam
 E_s : soil elastic modulus
 $e_1(x)$: virtual pressure for modeling shearing dislocation of joints
 $e_2(x)$: virtual pressure for modeling bending opening of joints
 F_1 : concentrated force exerting at joints
 GA : shear stiffness of beam
 $H(x-x_j)$: Heaviside step function
 I : second moment of inertia
 i : settlement trough width
 k : coefficient of subgrade reaction
 L : overall length
 l : length of difference nodes
 l_s : width of each segmental ring
 M_2 : a pair of couples applied at joints
 $M(x)$: bending moment of beam
 $M_e(x)$: virtual bending moment induced by virtual pressures
 $M_j(x)$: bending moment of longitudinal joint
 $Q(x)$: shear force of beam
 $Q_e(x)$: virtual shear force induced by virtual pressures
 $Q_j(x)$: shear force of longitudinal joint
 $q(x)$: additional pressure exerting on Timoshenko beam
 r : radius of tunnel
 $s(x)$: tunneling-induced subsurface settlement
 s_{\max} : maximum subsurface settlement
 V_1 : new tunneling-induced volume loss
 $w(x)$: deflection of Timoshenko beam
 x_j : location of longitudinal joint
 Z_0 : buried depth of new tunnel
 z_p : buried depth of pipeline
 α : coefficient of shearing stiffness reduction
 β : coefficient of bending stiffness reduction
 $\Delta\theta$: relative rotation angle
 $\Delta\delta$: relative dislocation
 $\delta'(x-x_j)$: first-order derivative of Dirac delta function
 $\delta''(x-x_j)$: second-order derivative of Dirac delta function

ψ : neutral-axis angle

μ : soil Poisson's ratio

Appendix B. Detailed expressions of equations, matrices, and vectors

The finite differential forms of virtual pressures $e_1(x)$ and $e_2(x)$ are as follows.

$$e_1(x_j) = \begin{cases} \frac{(1-\alpha)}{l} EI \left(\frac{d^3 w(x_j)}{dx^3} - \frac{k}{GA} \frac{dw(x_j)}{dx} + \frac{1}{GA} \frac{dq(x_j)}{dx} \right), & (i = j-1) \\ 0 & (i = j) \\ -\frac{(1-\alpha)}{l} GA \left(\frac{d^3 w(x_j)}{dx^3} - \frac{k}{GA} \frac{dw(x_j)}{dx} + \frac{1}{GA} \frac{dq(x_j)}{dx} \right), & (i = j+1) \end{cases} \quad (\text{B.1})$$

$$e_2(x_j) = \begin{cases} \frac{(1-\beta)}{l^2} EI \left(\frac{d^2 w(x_j)}{dx^2} - \frac{k}{GA} w(x_j) + \frac{q(x_j)}{GA} \right), & (i = j-1) \\ -2 \frac{(1-\beta)}{l^2} EI \left(\frac{d^2 w(x_j)}{dx^2} - \frac{k}{GA} w(x_j) + \frac{q(x_j)}{GA} \right), & (i = j) \\ \frac{(1-\beta)}{l^2} EI \left(\frac{d^2 w(x_j)}{dx^2} - \frac{k}{GA} w(x_j) + \frac{q(x_j)}{GA} \right), & (i = j+1) \end{cases}$$

The finite differential form of the governing differential equation is as follows.

$$\left. \begin{aligned}
& \frac{6w_i - 4(w_{i+1} + w_{i-1}) + (w_{i+2} + w_{i-2})}{l^4} - \frac{k}{GA} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} + \frac{k}{EI} w_i & (i \neq j-1, j, j+1) \\
& = \frac{1}{EI} q_i - \frac{1}{GA} \frac{q_{i+1} - 2q_i + q_{i-1}}{l^2}, \\
& \frac{6w_i - 4(w_{i+1} + w_{i-1}) + (w_{i+2} + w_{i-2})}{l^4} - \frac{1-\alpha}{l} \frac{-w_{i-2} + 2w_{i-1} - 2w_{i+1} + w_{i+2}}{2l^3} \\
& - \frac{k}{GA} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} - \frac{(1-\beta)}{l^2} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} + \frac{k}{GA} \frac{1-\alpha}{l} \frac{w_{i+1} - w_{i-1}}{2l} & (i = j-1) \\
& + \frac{k}{EI} w_i + \frac{k(1-\beta)}{GA l^2} w_i = \frac{1}{EI} q_i + \frac{(1-\beta)}{GA l^2} q_i + \frac{1-\alpha}{GA l} \frac{q_{i+1} - q_{i-1}}{2l} - \frac{1}{GA} \frac{q_{i+1} - 2q_i + q_{i-1}}{l^2}, \\
& \frac{6w_i - 4(w_{i+1} + w_{i-1}) + (w_{i+2} + w_{i-2})}{l^4} - \frac{k}{GA} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} + \frac{2(1-\beta)}{l^2} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} & (i = j) \\
& + \frac{k}{EI} w_i - \frac{2k(1-\beta)}{GA l^2} w_i = \frac{1}{EI} q_i - \frac{2(1-\beta)}{GA l^2} q_i - \frac{1}{GA} \frac{q_{i+1} - 2q_i + q_{i-1}}{l^2}, \\
& \frac{6w_i - 4(w_{i+1} + w_{i-1}) + (w_{i+2} + w_{i-2})}{l^4} + \frac{1-\alpha}{l} \frac{-w_{i-2} + 2w_{i-1} - 2w_{i+1} + w_{i+2}}{2l^3} \\
& - \frac{k}{GA} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} - \frac{(1-\beta)}{l^2} \frac{w_{i+1} - 2w_i + w_{i-1}}{l^2} - \frac{k}{GA} \frac{1-\alpha}{l} \frac{w_{i+1} - w_{i-1}}{2l} & (i = j+1) \\
& + \frac{k}{EI} w_i + \frac{k(1-\beta)}{GA l^2} w_i = \frac{1}{EI} q_i + \frac{(1-\beta)}{GA l^2} q_i - \frac{1-\alpha}{GA l} \frac{q_{i+1} - q_{i-1}}{2l} - \frac{1}{GA} \frac{q_{i+1} - 2q_i + q_{i-1}}{l^2}.
\end{aligned} \right\} \tag{B.2}$$

No virtual pressure, $e_1(x)$ or $e_2(x)$, was observed at the two ends of the shield tunnel. The four virtual nodes (w_{-2} , w_{-1} , w_{n+1} , and w_{n+2}) at each end of the tunnel were obtained because the boundary condition ensured that both the bending moments and shear forces were zero.

$$\begin{aligned}
w_{-1} &= \left(2 + \frac{kl^2}{GA}\right) w_0 - w_1 + \frac{l^2}{GA} q_0, \\
w_{-2} &= \left(2 + \frac{kl^2}{GA}\right) w_0 - 2\left(2 + \frac{kl^2}{GA}\right) w_1 + w_2 + \left(2 + \frac{kl^2}{GA}\right) \frac{l^2}{GA} q_0 + \frac{l^2}{GA} (q_1 - q_{-1}), \\
w_{n+1} &= \left(2 + \frac{kl^2}{GA}\right) w_n - w_{n-1} + \frac{l^2}{GA} q_n, \\
w_{n+2} &= \left(2 + \frac{kl^2}{GA}\right) w_n - 2\left(2 + \frac{kl^2}{GA}\right) w_{n-1} + w_{n-2} + \left(2 + \frac{kl^2}{GA}\right) \frac{l^2}{GA} q_n + \frac{l^2}{GA} (q_{n-1} - q_{n+1}).
\end{aligned} \tag{B.3}$$

The matrices and vectors in Eq. (20) are presented below.

$$\{Q_1\} = \begin{bmatrix} q_0 \\ q_1 \\ q_2 \\ \vdots \\ q_i \\ \vdots \\ q_{n-2} \\ q_{n-1} \\ q_n \end{bmatrix}_{1 \times (n+1)}, \quad (\text{B.16})$$

$$\{Q_2\} = \begin{bmatrix} q_1 - q_{-1} \\ q_2 - q_0 \\ q_3 - q_1 \\ \vdots \\ q_{i+1} - q_{i-1} \\ \vdots \\ q_{n-1} - q_{n-3} \\ q_n - q_{n-2} \\ q_{n+1} - q_{n-1} \end{bmatrix}_{1 \times (n+1)}, \quad (\text{B.17})$$

$$\{Q_3\} = \begin{bmatrix} q_1 - 2q_0 + q_{-1} \\ q_2 - 2q_1 + q_0 \\ q_3 - 2q_2 + q_1 \\ \vdots \\ q_{i+1} - 2q_i + q_{i-1} \\ \vdots \\ q_{n-1} - 2q_{n-2} + q_{n-3} \\ q_n - 2q_{n-1} + q_{n-2} \\ q_{n+1} - 2q_n + q_{n-1} \end{bmatrix}_{1 \times (n+1)}, \quad (\text{B.18})$$

$$\{Q_4\} = \begin{bmatrix} C_1 \\ C_2 \\ 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ C_3 \\ C_4 \end{bmatrix}_{1 \times (n+1)}, \quad (\text{B.19})$$

where

$$\begin{aligned}
A_1 &= \frac{kl^4 + 2(GA)^2}{(GA)^2}, A_2 = \frac{-2kl^2GA - 4(GA)^2}{(GA)^2}, A_3 = \frac{kl^2 - 2GA}{GA}, \\
B &= \frac{kl^2}{GA}, \\
C_1 &= \frac{GA(l^2q_0 - 2q_0 + q_1 - q_{-1}) + kl^2q_0}{l^2(GA)^2}, C_2 = \frac{q_0}{l^2GA}, \\
C_4 &= \frac{GA(l^2q_n - 2q_n + q_{n-1} - q_{n+1}) + kl^2q_n}{l^2(GA)^2}, C_3 = \frac{q_n}{l^2GA}, \tag{B.20}
\end{aligned}$$

$$\begin{aligned}
[D_{11}] &= \begin{bmatrix} -1 & 2 & 0 & -2 & 1 \\ 0 & 0 & 0 & 0 & 0 \\ 1 & -2 & 0 & 2 & -1 \end{bmatrix}_{3 \times 5}, [D_{12}] = \begin{bmatrix} -1 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & -1 \end{bmatrix}_{3 \times 3}, [D_{13}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}_{3 \times 3} \\
[D_{21}] &= \begin{bmatrix} 1 & -2 & 1 \\ -2 & 4 & -2 \\ 1 & -2 & 1 \end{bmatrix}_{3 \times 3}, [D_{22}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & -2 & 0 \\ 0 & 0 & 1 \end{bmatrix}_{3 \times 3},
\end{aligned}$$

where $[D_{11}]$, $[D_{12}]$, $[D_{13}]$, $[D_{21}]$, and $[D_{22}]$ are the stiffness matrices of the longitudinal joints. Considering an existing structure with multiple longitudinal joints, $[E_{11}]$, $[E_{12}]$, $[QE_1]$, $[E_{21}]$, $[E_{22}]$, and $[QE_2]$ will feature several $[D_{11}]$, $[D_{12}]$, $[D_{13}]$, $[D_{21}]$, and $[D_{22}]$ at the positive diagonal lines, respectively.